

University of Vermont

UVM ScholarWorks

College of Engineering and Mathematical
Sciences Faculty Publications

College of Engineering and Mathematical
Sciences

4-8-2013

Limited imitation contagion on random networks: Chaos, universality, and unpredictability

Peter Sheridan Dodds
University of Vermont

Kameron Decker Harris
University of Vermont

Christopher M. Danforth
University of Vermont

Follow this and additional works at: <https://scholarworks.uvm.edu/cemsfac>



Part of the [Human Ecology Commons](#), and the [Medicine and Health Commons](#)

Recommended Citation

Dodds PS, Harris KD, Danforth CM. Limited imitation contagion on random networks: Chaos, universality, and unpredictability. *Physical review letters*. 2013 Apr 8;110(15):158701.

This Article is brought to you for free and open access by the College of Engineering and Mathematical Sciences at UVM ScholarWorks. It has been accepted for inclusion in College of Engineering and Mathematical Sciences Faculty Publications by an authorized administrator of UVM ScholarWorks. For more information, please contact scholarworks@uvm.edu.

Limited Imitation Contagion on Random Networks: Chaos, Universality, and Unpredictability

Peter Sheridan Dodds,^{*} Kameron Decker Harris,[†] and Christopher M. Danforth[‡]

Department of Mathematics and Statistics, Computational Story Lab, Vermont Complex Systems Center, and Vermont Advanced Computing Core, The University of Vermont, Burlington, Vermont 05401, USA

(Received 1 August 2012; revised manuscript received 7 March 2013; published 8 April 2013)

We study a family of binary state, socially inspired contagion models which incorporate imitation limited by an aversion to complete conformity. We uncover rich behavior in our models whether operating with either probabilistic or deterministic individual response functions on both dynamic and fixed random networks. In particular, we find significant variation in the limiting behavior of a population's infected fraction, ranging from steady state to chaotic. We show that period doubling arises as we increase the average node degree, and that the universality class of this well-known route to chaos depends on the interaction structure of random networks rather than the microscopic behavior of individual nodes. We find that increasing the fixedness of the system tends to stabilize the infected fraction, yet disjoint, multiple equilibria are possible depending solely on the choice of the initially infected node.

DOI: [10.1103/PhysRevLett.110.158701](https://doi.org/10.1103/PhysRevLett.110.158701)

PACS numbers: 89.65.-s, 05.45.-a, 87.23.Ge

The structure and dynamics of real, complex networks remains an open area of great research interest, particularly in the realm of evolutionary processes acting on and within networked systems [1–6]. Here, motivated by considerations of social contagion—the spreading of ideas and behaviors between people through social networks and media—we explore an idealized, binary-state social contagion model in which individuals choose to be like others but only up to a point: they do not want to be like everyone else [7–10]. We term such behavior “limited imitation contagion.” We build naturally on previous studies of threshold models of contagion [12–15], and our model can also be seen as a specific subfamily of dynamical Boolean network models [16,17]. We show how macroscopic network structure overrides microscopic details, and we find complex dynamics whose character moves from universal and predictable to particular and unpredictable as we allow the system to become increasingly deterministic.

In constructing our model, our main interest is in understanding how spreading by limited imitation contagion on random networks behaves under three main tunable conditions: (1) social awareness: the rate of contact between individuals; (2) social variability: the extent to which friendships are fixed; and (3) social influence: the character of individuals' responses to the behavior of others.

To begin with, we consider a binary state model for which individuals are either in a base state S_0 or an alternate state S_1 . We assume individuals interact over an uncorrelated random network, which may be dynamic or fixed. For simplicity, and due to the richness of the dynamics we find, we employ standard Erdős-Rényi networks which possess Poisson degree distributions. We take time to be discrete ($t = 0, 1, 2, \dots$), and we prescribe each node's degree k at $t = 0$. In a dynamic network, when node i updates, it samples the states of k_i randomly chosen nodes (i.e., the system is a random mixing model

with nonuniform contact rates). For a fixed network, node i repeatedly samples the same k_i nodes. We further restrict our attention to single-seed contagion processes wherein all nodes are in state S_0 at time $t = 0$, with one randomly chosen node in state S_1 .

The contagion process is manifested through the response functions of individual nodes. We allow nodes to update synchronously, and node i 's response function $F_i: [0, 1] \mapsto [0, 1]$ gives the probability that node i will be in state S_1 upon updating, where the argument taken by F_i is the fraction of nodes sampled by node i that are currently in state S_1 , $\phi_{i,t}$.

We investigate two kinds of response functions, probabilistic and deterministic, both of which incorporate the

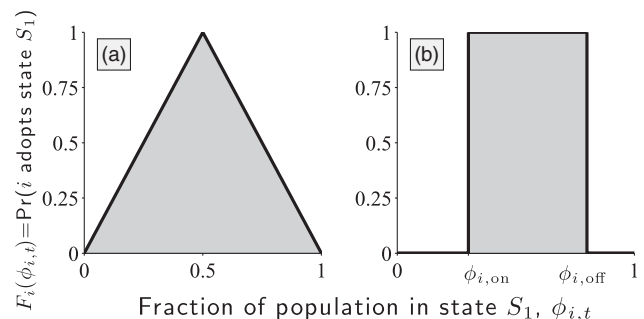


FIG. 1. Examples of probabilistic and deterministic response functions capturing limited imitation contagion dynamics. At each time t , nodes use their given response functions to update their own state based on the perceived fraction of their neighbors in state S_1 , $\phi_{i,t}$. We construct the tent map T_2 (see main text) shown in (a) by averaging over deterministic response functions of the kind shown in (b) by considering a family of the latter with “on” and “off” thresholds uniformly distributed in $[0, \frac{1}{2}]$ and $[\frac{1}{2}, 1]$, respectively. We then build networked systems whose macroscopic character is tent map-like but differ strongly at the microscopic level.

characteristic of the adoption probability growing and then diminishing as the perceived popularity of S_1 increases. In Fig. 1(a), we show an example of a probabilistic response function, the tent map, which is defined as $T_r(x) = rx$ for $0 \leq x \leq \frac{1}{2}$ and $r(1-x)$ for $\frac{1}{2} \leq x \leq 1$. Here, we consider the $r = 2$ case, meaning node i adopts state S_1 with probability $T_2(\phi_{i,t})$. We use the tent map T_2 for a number of reasons. (1) As a standard iterative map of the unit interval, the tent map’s dynamics for $r = 2$ are both interesting and well understood [18]—it is fully chaotic and its invariant density χ is uniform on $[0,1]$ (χ is the long-term probability distribution for the values of a map’s iterates). (2) The tent map captures a probabilistic flavor of the adopt-when-novel, drop-when-ubiquitous behavior we aim to model. (3) We can construct a simple and elegant connection with deterministic contagion processes which we describe next.

In Fig. 1(b), we show an example of a deterministic response function which is characterized by “on” and “off” thresholds, ϕ_{on} and ϕ_{off} . Node i will only adopt or remain in state S_1 if it perceives the fraction of others in state S_1 to lie between its on and off thresholds: $\phi_{i,on} \leq \phi_{i,t} < \phi_{i,off}$.

For these deterministic “on-off” response functions, we examine the special case where ϕ_{on} is distributed uniformly

on $[0, \frac{1}{2}]$ and ϕ_{off} likewise on $[\frac{1}{2}, 1]$. Averaging over all deterministic response functions created in this way, we obtain precisely the tent map T_2 . Collectively, we then have the same on-average behavior in both the probabilistic and deterministic cases, and this allows us to profitably explore the effect of varying the specific response dynamics at the micro level, as well as interaction patterns.

We now examine the behavior of four subfamilies of our model that vary in terms of node response functions being probabilistic or deterministic (P or D), and whether or not network connections randomly rewire or are fixed (R or F). We refer to these model classes as P - R , P - F , D - R , and D - F . Each class is indexed by the parameter of average node degree k_{av} , and, as per our design, all four systems have the same on-average response functions (the tent map T_2) and degree distributions (Poisson).

Our first focus is on the long-term behavior of these four networked systems. Key to our understanding are the well-known behaviors of the tent and logistic maps acting iteratively on the unit interval, the former given above and the latter by $L_r(x) = rx(1-x)$. We reproduce the orbit diagrams for these models in Figs. 2(a) and 2(b). Both systems are controlled by the amplitude parameter r , whose increase leads to changes in their invariant densities,

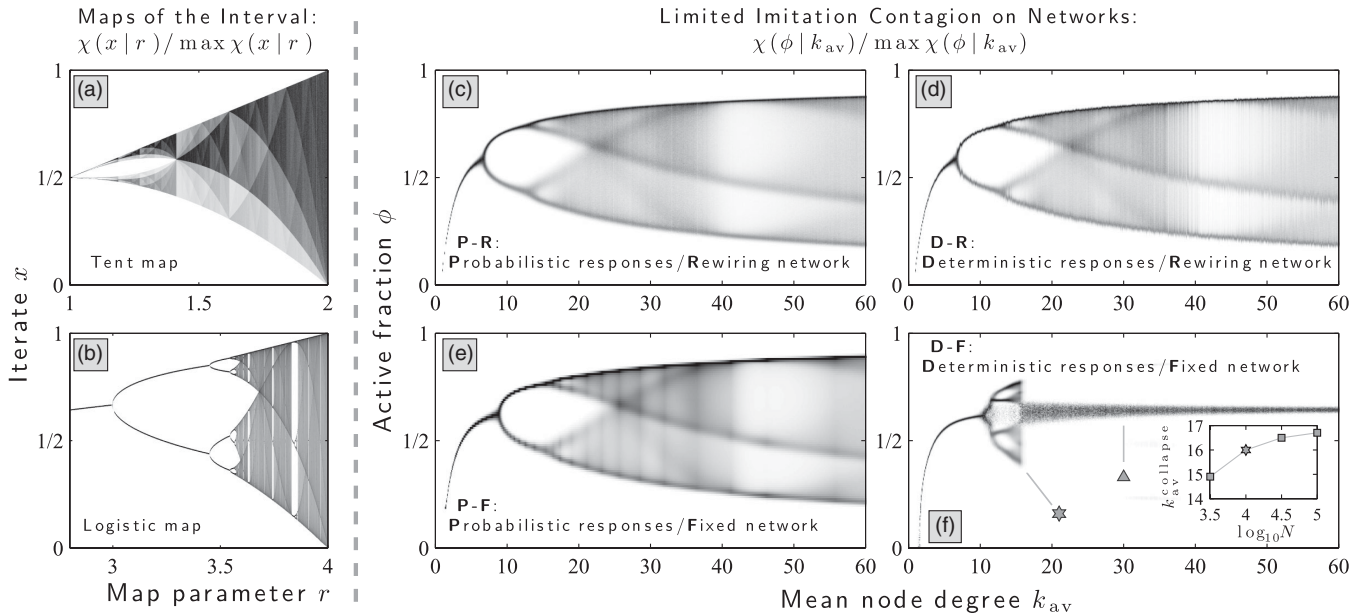


FIG. 2. Comparison of bifurcation diagrams for networked systems with limited imitation contagion based on tent map response functions. (a), (b) For reference, standard bifurcation diagrams for the tent map and the logistic map operating as maps of the unit interval, showing distinct universality classes. The gray scale represents the normalized invariant density χ (darker indicating higher values). (c)–(f) Orbit diagrams for limited imitation contagion acting on standard random networks for four model classes with probabilistic or deterministic responses and network interactions as indicated. The tunable parameter is average degree k_{av} . The inset in (f) shows the approximate location of the collapse point $k_{av}^{collapse}$ as a function of network size, the stars indicating the collapse point for $N = 10^4$. The triangle in (f) indicates $k_{av} = 30$ for the system examined in Fig. 3. Simulation details: network size $N = 10^4$; one random seed per network; granularity of 0.1 in k_{av} ; binning size of 0.001 in ϕ ; for P - R , D - R , and P - F systems, we ran simulations for 10^4 time steps on 10^3 networks, taking values of ϕ_t for $t \geq 10^3$; and for D - F systems, which exhibit long transient behavior (see text), we ran simulations for 10^5 time steps on 200 networks, ignoring the first 5×10^4 time steps.

famously resulting in the bifurcation diagram for the logistic map. The two maps, while topologically conjugate, produce distinct bifurcation diagrams.

As we show in Figs. 2(c)–2(f), the signature of the microscopic response function of the tent map is erased by the network dynamics of the four model classes. Macroscopically, we see four orbit diagrams analogous to the logistic map’s characteristic bifurcation diagram. We see that increasing the average connectivity of the network k_{av} is equivalent to increasing the logistic map’s amplitude parameter r , and the system moves along the period-doubling route to chaos.

However, while appearing to belong to the same universality class, the orbit diagrams of the four models differ importantly in detail, most profoundly for the fully deterministic $D-F$ class.

First, we observe that the four systems produce only three distinct orbit diagrams, since for large enough systems, the two random mixing classes $P-R$ and $D-R$ must exhibit the same macroscopic behavior [Figs. 2(c) and 2(d)]. For the $D-R$ class, random rewiring overwhelms the fact of each node having a fixed deterministic response function [for finite systems some evidence of discreteness limits the smoothness of the orbit diagram in Fig. 2(d)]. While we present simulation results only here, we note that we are able to fully address the $P-R$ and $D-R$ models analytically, and we find excellent agreement (see Ref. [19] where we explore more general influence maps).

In next considering the fixed network model with probabilistic response functions, $P-R$, we see in Fig. 2(e) that the orbit diagram has slightly moved to the right—the bifurcation points now occur for higher values of k_{av} relative to the random mixing cases. A robust, discernible vertical striation also appears, with separation between stripes increasing with k_{av} , for which we do not have an explanation. Thus, making the network fixed appears to induce some modest changes in the dynamics, though we cannot discount finite size effects.

Our final model class, $D-F$ —to whose description we devote the remainder of the Letter—is the most structured, and it generates considerably different and intricate behavior. Once the network and set of response functions is realized, the system is now completely deterministic (and irreversible), and we find surprising changes in the system’s behavior at both macroscopic and microscopic scales. In Fig. 2(f), we see that the orbit diagram for the $D-F$ model collapses abruptly after several rounds of period doubling, which are themselves relatively compressed in terms of k_{av} . Our simulations suggest that above an average degree of $k_{av} \approx 17$, the macroscopic dynamics always collapse. The inset in Fig. 2(f) shows estimates of the collapse point for $N = 10^{3.5}$ up to $N = 10^5$, and the asymptotic growth with $\log_{10} N$ strongly suggests the collapse is real and not a finite size effect. The collapse appears to favor a fixed point macroscopic state, around $\phi = 2/3$, which is the fixed point of the tent map T_2 .

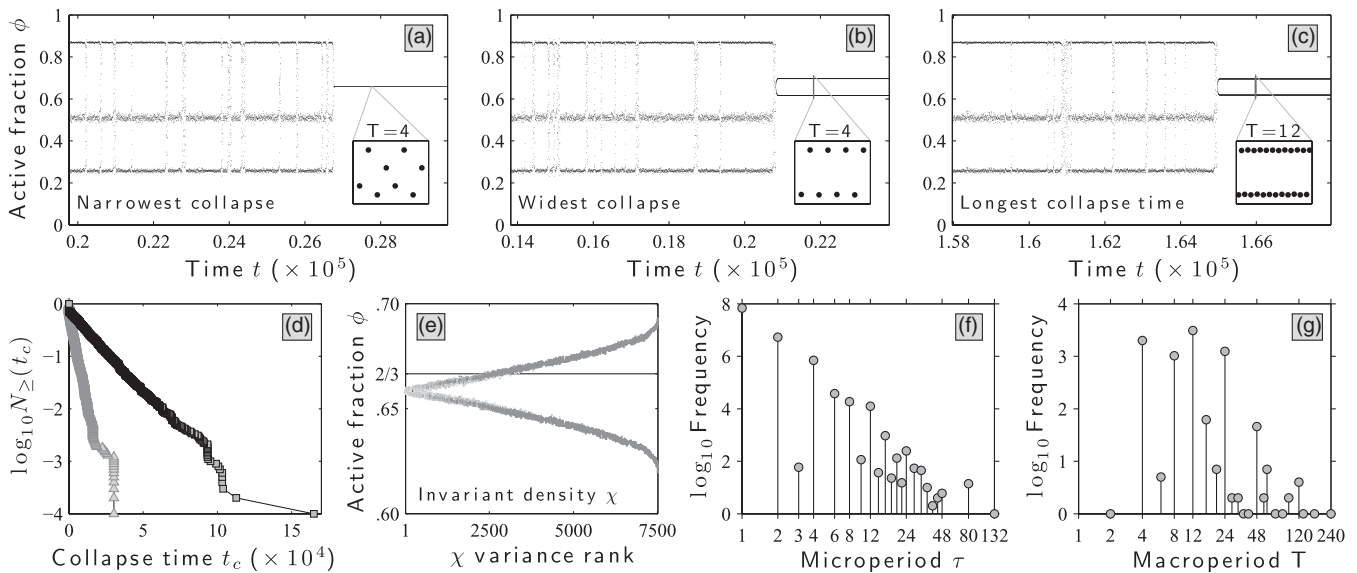


FIG. 3. Comprehensive survey of the possible limited imitation contagion dynamics for a single fixed random network ($k_{av} = 30$) with its $N = 10^4$ nodes having fixed deterministic response functions. We ran N simulations, activating each node as a single seed and allowed the fully deterministic dynamics to run until collapse. (a)–(c) Partial time series for three different initial seeds leading to the narrowest and widest postcollapse dynamics (a),(b) and the longest collapse time t_c (c). The insets in (a)–(c) magnify two cycles of the underlying periodic behavior. (d) Complementary cumulative distribution of collapse times (black squares) compared with an example distribution from another network (gray triangles). (e) Invariant densities, ordered by variance. (f) Histogram of postcollapse microperiod τ . (g) Histogram of postcollapse macroperiod T .

However, a closer examination of the D - F class's potential dynamics reveals a far more subtle story.

In Fig. 3, we summarize the possible dynamics of a lone realized network ($N = 10^4$, $k_{\text{av}} = 30$) with a single set of fixed deterministic response functions. We exhaustively ran $N = 10^4$ tests of the system's behavior by separately seeding each individual node. Of these, 7515 contagion events successfully lead to long-term, nonzero infection levels. Figures 3(a)–3(c) show partial time series for three system evolutions for the same network, differing only by seed node, and corresponding respectively to the narrowest and widest collapses, as measured by the variance of χ , and the longest time to collapse. In each case, the initial dynamics of ϕ exhibit a clear period three pattern interspersed with intermittent chaotic dynamics, followed by a sharp collapse to distinct periodic behaviors.

Across all initial seeds, the time to collapse t_c varies greatly, with the extreme example of Fig. 3(c) collapsing after more than 160 000 time steps and over 1.6×10^9 individual node updates. In Fig. 3(d), we provide the complementary cumulative distribution for collapse times (black squares). The semilog scale indicates that an exponential decay for $t_c < 10^5$ covers the majority of cases with the longest collapse time clearly a singular outlier. Even after the dynamics collapse, we see that on the order of 10% of all nodes change state in each time step. Surprisingly, we see different distributions for t_c for different individual networks. For example, the gray triangles represent the collapse time distribution for another randomly realized network with $k_{\text{av}} = 30$. Again we see an exponential distribution but now with a much steeper decay. Rather than finding one outlier, we see that a set of exceptional dynamics lasts for much longer than the main exponential decay would suggest, with an isolated set distributed around $t_c \approx 25\,000$, and a group ending at around $t_c \approx 30\,000$.

We turn lastly to the behavior after the collapse. We call the postcollapse period of ϕ_t the system's "macroperiod" T , and the period of an individual node its "microperiod" τ . Figure 3(e) shows the invariant density χ for all seeds post collapse, ordered by χ 's variance (contagion events that fail are not included). Strikingly, different seeds almost always lead to distinct invariant densities, with an apparent dominance of period $T = 2$ behavior (but see below). We also find that the tent map's fixed point of $2/3$ [horizontal line in Fig. 3(e)] is not the center of the invariant densities.

In Figs. 3(f) and 3(g), we provide complete histograms of all postcollapse microperiods and macroperiods for this particular network (again ignoring failed spreading events). Unlike the collapse time distribution, we see broadly similar distributions for other individual networks. Microperiods such as 1, 2, 4, 12, and 24 form a dominant envelope, and many odd-numbered microperiods are absent. We find that the most common resultant macroperiods are 4, 8, 12, and 24, that a pure macroperiod 2 is relatively rare ($< 0.1\%$), and that the largest observed

macroperiod is 240. Again, all of these outcomes are deterministic, depending only on the choice of initial seed.

To conclude, in abstracting from a real world problem, we have constructed a rich, networked-based model of limited imitation contagion that allows us to deeply explore the effects of manifesting either deterministic or probabilistic microscopic behavior in both network structure and node response in a well-defined fashion. Our findings broadly suggest that increasing the degree of a social system's microscopic rigidity leads to higher levels of long-term unpredictability at macroscopic scales. Future work could take at least two paths, one directed toward the dynamical system's aspects and the other toward a more realistic empirically supported model. Some natural objectives would be to explain why the collapse occurs for the D - F class, to explore how local properties of seed nodes such as node degree, correlations, and clustering relate to the collapse time and postcollapse dynamics, to test other contagion mechanisms across systems of increasingly fixed microscopic structure, and to examine these limited imitation contagion models on real networks and classes of naturally occurring ones.

The authors thank A. Mandel and D. J. Watts for encouragement, and are grateful for the computational resources provided by the Vermont Advanced Computing Core supported by NASA (NNX 08A096G). K. D. H. was supported by VT-NASA EPSCoR; C. M. D. and P. S. D. were supported by a grant from the MITRE Corporation; P. S. D. was supported by NSF CAREER Grant No. 0846668.

*peter.dodds@uvm.edu

†kameron.harris@uvm.edu

‡chris.danforth@uvm.edu

- [1] M. E. J. Newman, *SIAM Rev.* **45**, 167 (2003).
- [2] D. J. Watts and P. S. Dodds, *J. Consum. Res.* **34**, 441 (2007).
- [3] V. Colizza, A. Barrat, M. Barthelmeij, A.-J. Valleron, and A. Vespignani, *PLoS Med.* **4**, e13 (2007).
- [4] S. Wuchty, B. F. Jones, and B. Uzzi, *Science* **316**, 1036 (2007).
- [5] C. A. Hidalgo, B. Klinger, A.-L. Barabási, and R. Hausman, *Science* **317**, 482 (2007).
- [6] Y.-Y. Liu, J.-J. Slotine, and A.-L. Barabási, *Nature (London)* **473**, 167 (2011).
- [7] G. Simmel, *Am. J. Sociology* **62**, 541 (1957).
- [8] M. S. Granovetter and R. Soong, *J. Econ. Behav. Organ.* **7**, 83 (1986).
- [9] J. Ugander, L. Backstrom, C. Marlow, and J. Kleinberg, *Proc. Natl. Acad. Sci. U.S.A.* **109**, 5962 (2012).
- [10] Our model is distinct from traditional coupled maps [11], having binary rather than continuous node states, and exhibits fundamentally different behavior.
- [11] K. Kaneko, *Prog. Theor. Phys.* **72**, 480 (1984).
- [12] T. C. Schelling, *J. Math Sociology* **1**, 143 (1971).
- [13] M. Granovetter, *Am. J. Sociology* **83**, 1420 (1978).
- [14] D. J. Watts, *Proc. Natl. Acad. Sci. U.S.A.* **99**, 5766 (2002).
- [15] J. P. Gleeson, *Phys. Rev. E* **77**, 046117 (2008).

-
- [16] S. A. Kauffman, *J. Theor. Biol.* **22**, 437 (1969).
- [17] M. Aldana, S. Coppersmith, and L. P. Kadanoff, in *Perspectives and Problems in Nonlinear Science*, edited by E. Kaplan, J. E. Marsden, and K. R. Sreenivasan (Springer, New York, 2003), Chap. 2, pp. 23–90.
- [18] K. T. Alligood, T. D. Sauer, and J. A. Yorke, *Chaos: An Introduction to Dynamical Systems* (Springer, New York, 1996).
- [19] K. D. Harris, C. M. Danforth, and P. S. Dodds, [arXiv:1303.1414](https://arxiv.org/abs/1303.1414).